

AUTOMATING SYMMETRY-BREAKING CALCULATIONS

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Abstract

The process of classifying possible symmetry-breaking bifurcations requires a computation involving the subgroups and irreducible representations of the original symmetry group. It is shown how this calculation can be automated using a group theory package such as GAP. This enables a number of new results to be obtained for larger symmetry groups, where manual computation is impractical. Examples of symmetric and alternating groups are given, and the method is also applied to the spatial symmetry-breaking of periodic patterns observed in experiments.

1. *Introduction*

This paper is concerned with symmetry-breaking bifurcations, in which a state with a high degree of symmetry becomes unstable to a state with less symmetry. As an introductory example, consider an object with symmetry (for example, spherical, cubic or square symmetry) subjected to steadily increasing external pressure. As the pressure increases, the object may retain its symmetry, or, at a critical value of the pressure, it may spontaneously buckle into a different shape with less symmetry. More generally, symmetry-breaking bifurcations occur in systems of ordinary or partial differential equations with a certain symmetry; there is a solution of the equations with this full symmetry, but this state may become unstable as a parameter is varied. Because of the symmetry, the linear stability problem has multiple eigenvalues, with a multiplicity given by the dimension of one of the irreducible representations of the symmetry group. This means that the linear stability problem is degenerate, and this degeneracy can be resolved only by considering the nonlinear terms in the equations, which are themselves constrained by the symmetry. The general theory for such transitions is set out in the key text on the subject by Golubitsky, Stewart and Schaeffer [11], building on earlier work by Sattinger [22]; two more recent books incorporate newer theoretical results and provide many examples and applications of the theory [4, 10].

A remarkable feature of such problems is that the possible symmetry-breaking transitions can be determined purely from the symmetry of the problem, independently of the underlying physical mechanisms involved, by the use of group and representation theory. However, in order to determine which of the possible transitions lead(s) to a stable solution in a particular application, it is necessary to return to the original equations and compute the values of certain coefficients.

One application of considerable interest is the formation and stability of patterns. In experiments on Rayleigh–Bénard convection, where a layer of fluid is heated from below, and in the Faraday experiment, where a layer of fluid or granular material is subjected to a vertical oscillation, the original Euclidean symmetry can be broken, leading to a regular

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pattern of stripes, squares or hexagons [21, 3, 1]. These regular periodic patterns may themselves become unstable, leading to secondary patterns that may retain periodicity on a lengthscale longer than that of the original pattern [14, 24, 6, 7].

The aim of this paper is to show how the basic group-theoretic approach to classifying possible symmetry-breaking bifurcations can be straightforwardly automated by using a computational group theory package such as GAP [8]. Section 2 summarises the general theory of the subject, and Section 3 describes the computational implementation. A number of results are presented in tabular form in Section 4, and these are extended to the case of oscillatory bifurcations in Section 5.

2. Problem formulation

This section introduces the language and notation of bifurcation with symmetry, and summarises the key results; full details can be found in [11] or [4]. Suppose that G is a compact Lie group, acting absolutely irreducibly on $V = \mathbb{R}^n$. The representation is said to be *absolutely irreducible* if the only matrices that commute with all elements of G are multiples of the identity, and this implies that the representation is irreducible (that is, there is no proper G -invariant subspace).

The system of n equations

$$f(x, \lambda) = 0, \quad x \in \mathbb{R}^n, \lambda \in \mathbb{R} \quad (1)$$

is said to be G -equivariant if

$$f(gx, \lambda) = gf(x, \lambda) \quad \text{for all } g \in G, \quad (2)$$

where the notation gx is used as an abbreviation for the action of g on x through the matrix representation of g .

The symmetry of a point v in V is described by its isotropy subgroup, or stabiliser, defined by

$$\{g \in G : gv = v\}.$$

A subgroup $H \subset G$ is said to be an *isotropy subgroup* of G if there exists a point $v \in V$ with isotropy H . In the physics literature, isotropy subgroups are often referred to as *little groups*. For any subgroup $H \subset G$, the fixed-point subspace of H is

$$\text{Fix}(H) = \{v \in V : hv = v \text{ for all } h \in H\}. \quad (3)$$

The irreducibility condition means that either $\text{Fix}(G) = V$ or $\text{Fix}(G) = 0$. However, if $\text{Fix}(G) = V$, then any subspace of V is G -invariant, which contradicts the irreducibility if $n > 1$. Hence for $n > 1$, we must have $\text{Fix}(G) = 0$.

We now suppose that the system of equations (1) has a solution that undergoes a stationary bifurcation at $x = 0, \lambda = 0$, so that

$$f(0, 0) = 0 \quad \text{and} \quad \left| \frac{\partial f_i}{\partial x_j} \right| (0, 0) = 0. \quad (4)$$

The second condition, that the Jacobian matrix has zero determinant, means that the implicit function theorem does not hold, and so (1) does not define a unique branch of solutions $x(\lambda)$ in the vicinity of $\lambda = 0$. The linearisation of f must be a scalar multiple $c(\lambda)$ of the identity (from absolute irreducibility) with $c(0) = 0$, and it is assumed that the transversality condition $c'(0) \neq 0$ holds.

The type of the bifurcation depends on the choice of representation of G . For the trivial representation, in which $gx = x$ for all g and x , the equivariance condition (2) provides no constraint, and generically the bifurcation is of saddle-node type, in which the leading terms are, after a suitable rescaling, $\lambda + x^2 = 0$. There is no loss of symmetry at the bifurcation. For all other one-dimensional representations, $gx = -x$ for some g , and so (2) forces f to be an odd function of x , and the bifurcation is of pitchfork type, $\lambda x + ax^3 = 0$ to leading order. The nonzero solutions have an isotropy subgroup generated by the elements of G that act as $+1$.

For irreducible representations (irreps) of dimension $n > 1$, we must have $\text{Fix}(G) = 0$, as explained above. Applying (2) to $x = 0$ gives $f(0, \lambda) = gf(0, \lambda)$ for all $g \in G$, so that $f(0, \lambda) \in \text{Fix}(G)$, and hence $f(0, \lambda) = 0$. Physically, this is the statement that a solution with the full symmetry group G exists for all values of the parameter λ . The existence of non-zero solutions in this case can be established using the equivariant branching lemma [5, 25, 13, 11], which states that if, in addition to the conditions given above, there is an isotropy subgroup H of G with $\text{Dim}(\text{Fix}(H)) = 1$, then a unique branch of solutions to (1) with isotropy H exists, emanating from $x = 0, \lambda = 0$. This condition is sufficient but not necessary, so applying the lemma in general finds some but not all of the solution branches of (1). The notation $D(H)$ will be used henceforth as an abbreviation for $\text{Dim}(\text{Fix}(H))$. If $D(H) > 1$, there is no simple analogue of the equivariant branching lemma; however, if $D(H) = 2$, then at least one solution exists in $\text{Fix}(H)$, provided that the quadratic terms in the equations are non-zero and satisfy certain non-degeneracy conditions [15, 16, 18].

In the case where $n > 1$, the type of bifurcation is constrained by the normalizer, $N(H)$, defined by

$$N(H) = \{g \in G : g^{-1}Hg = H\}. \tag{5}$$

Equivalently, this is the largest subgroup of G in which H is normal. The normalizer maps $\text{Fix}(H)$ to itself, since if $y \in \text{Fix}(H)$, $g \in N(H)$ and $h \in H$, then $(g^{-1}hg)y = y$ (since $g^{-1}hg \in H$), so $hgy = gy$ and hence $gy \in \text{Fix}(H)$. Clearly, $N(H)$ contains H , and if $N(H) = H$ then there is no symmetry within $\text{Fix}(H)$, since all elements of H act as the identity in $\text{Fix}(H)$. But if $N(H) \neq H$, then the symmetry group acting on $\text{Fix}(H)$ is the quotient group $N(H)/H$. If $D(H) = 1$, the only possibilities are either that $N(H) = H$, or that $N(H)/H$ is isomorphic to the cyclic group Z_2 ; in the latter case, the bifurcation is of pitchfork type. If $N(H) = H$, then the bifurcation is generally transcritical, but in some cases there may be additional constraints, for example due to ‘hidden’ symmetries [10, 7]. A standard example is the group D_5 , where $N(H) = H$ but the bifurcation is not transcritical.

In order to apply the equivariant branching lemma and hence obtain solutions of (1), one must find the isotropy subgroups H for which $D(H) = 1$. These are often referred to as *axial* subgroups, since they correspond to axes of reflection for two-dimensional irreps, and to axes of rotation for three-dimensional irreps (for the cube, for example, there are three such axes, with isotropies Z_4, Z_3 and Z_2). To compute $D(H)$, it is not necessary to have an explicit form of the matrix representation, since one can use the trace formula [11, 4]

$$D(H) = \text{Dim}(\text{Fix}(H)) = \frac{1}{|H|} \sum_{h \in H} \chi(h), \tag{6}$$

where $\chi(h)$ is the character of h (the trace of the matrix). Thus the computation of $D(H)$ requires only the character table of G .

3. Computational algorithm

The algorithm for computing the isotropy subgroups of a given finite group G is outlined in Table 1. This algorithm was implemented using the computational group theory package GAP [8].

The first step is to define the group G , by specifying the generators of G . In most cases, the most convenient way to do this is to write G as a subgroup of a permutation group, which can always be done (this is Cayley's theorem). For example, the group D_4 of symmetries of a square is expressed as a subgroup of S_4 by labelling the four corners; the two reflections that generate D_4 are then the permutations $(1, 2)(3, 4)$ and $(1, 3)$. The next stage is to find the lattice of subgroups of G , up to conjugacy. Two subgroups H and K are said to be *conjugate* if there exists $g \in G$ such that $H = g^{-1}Kg$; for example, the subgroups generated by $(1, 3)$ and $(2, 4)$ are conjugate, both corresponding to diagonal reflections of the square. If two subgroups are conjugate, then the dimensions of their fixed-point subspaces are equal. Note that two subgroups may be isomorphic but not conjugate; for example, D_4 has three non-conjugate subgroups that are isomorphic to Z_2 . For each subgroup H , it is useful to check whether H is normal, and to find the normaliser of H .

The next stage is to find the character table of G . Note that in a high-level language such as GAP, this is done in a single command (as is the construction of the subgroup lattice). We now choose one of the irreps of G , and calculate $D(H)$ for all non-conjugate subgroups of G using the trace formula (6). In the GAP language, this is evaluated as a scalar product of the restriction of the character to H with the trivial character of H . Having found the dimensions corresponding to each subgroup, the isotropy subgroups can be found, using

Table 1: Computational algorithm for finding isotropy subgroups.

Specify the group G as a subgroup of a permutation group
Compute the lattice of non-conjugate subgroups of G
Compute the character table of G
For each irrep of G
For each subgroup H of G
Compute $D(H)$ using trace formula
End for
For each subgroup H of G
If $D(H) > 0$ and $D(K) < D(H)$ for all K immediately above H in the lattice, then H is an isotropy subgroup
If $D(H) > 0$ and $D(K) = 0$ for all K immediately above H in the lattice, then H is maximal
If H is an isotropy subgroup, find $N(H)/H$
End for
List isotropy subgroups H and $D(H)$ for this irrep
End for

the chain criterion (see, for example, [20]). This states that H is an isotropy subgroup of G if and only if $D(K) < D(H)$ for all $K \supsetneq H$. Since $K \supsetneq H \Rightarrow D(K) \leq D(H)$, it is only necessary to check whether $D(K) < D(H)$ for all K immediately above H in the subgroup lattice. Note that this chain criterion is valid only for finite groups; refinements have been proposed to provide a criterion for infinite groups, such as $SO(3)$ and $O(3)$; see [13, 17]. The group G is itself always an isotropy subgroup, with $D(G) = 0$ for all non-trivial irreps. Any subgroup $H \subsetneq G$ with $D(H) = 0$ cannot be an isotropy subgroup, by the chain criterion, so it is only necessary to find isotropy subgroups with $D(H) \geq 1$. Having found an isotropy subgroup, it is useful to compute the quotient group $N(H)/H$, since this gives the symmetry acting on $\text{Fix}(H)$, which influences the type of bifurcation.

Amongst isotropy subgroups, of particular interest are the maximal isotropy subgroups. An isotropy subgroup H is said to be *maximal* if there is no isotropy subgroup K satisfying $H \subsetneq K \subsetneq G$. Equivalently, H is maximal if $D(H) > 0$ and $D(K) = 0$ for all K immediately above H in the subgroup lattice. Thus isotropy subgroups with $D(H) = 1$ are always maximal, but those with $D(H) > 1$ may or may not be maximal. Having found the isotropy subgroups for a given irrep, the computation is then repeated for the other irreps of G .

There are a number of short cuts that can be taken to reduce the calculations required. These may be useful when the calculation is being carried out by hand, but it turns out that there is little to be gained from implementing these computationally, since the saving in computational time is not significant. The most time-consuming part of the computation is the construction of the subgroup lattice; applying the trace formula and the chain criterion typically occupies less than one fifth of the computation time.

The first short cut is that it is really only necessary to consider faithful irreps (recall that an irrep is *faithful* if only the identity element of g satisfies $gv = v$ for all $v \in V$). If the irrep is not faithful, the calculation can be reduced to the case of a smaller group, which is the quotient group G/N , where N is the normal subgroup of G containing the elements corresponding to the identity matrix. A standard example (see [11]) is the two-dimensional irrep of S_4 , in which a normal subgroup isomorphic to D_2 acts as the identity and the quotient group is isomorphic to S_3 . Faithful irreps can be identified from the character table from the fact that all (non-identity) conjugacy classes have a character less than that of the identity. For a faithful irrep of dimension n , it follows that the identity subgroup I is an isotropy subgroup with $D(I) = n$, while for a non-faithful irrep, the above normal subgroup N is an isotropy subgroup with $D(N) = n$.

The second short cut is that for faithful irreps, normal subgroups can never be isotropy subgroups. This simple result does not appear to be given explicitly in the standard literature, so it is formally stated and proved here.

LEMMA 1. *Suppose that G acts irreducibly and faithfully on \mathbb{R}^n , and that H is a proper normal subgroup of G . Then $D(H) = 0$, and so H is not an isotropy subgroup of G .*

Proof. If H is a normal subgroup of G , then the normaliser of H is $N(H) = G$, by definition. Then, as shown in Section 2, G maps $\text{Fix}(H)$ to itself, so the subspace $\text{Fix}(H)$ is G -invariant. Since the action of G is irreducible, this subspace must be either \mathbb{R}^n or 0. Now if $\text{Fix}(H) = \mathbb{R}^n$ and H is not the identity subgroup, then there are non-identity elements that act as the identity, so the action cannot be faithful. Hence if the action is faithful, $\text{Fix}(H) = \{0\}$, and so $D(H) = 0$ and H is not an isotropy subgroup. \square

Thirdly, it is in fact not necessary to evaluate $D(H)$ for every group in the lattice. For example, once a subgroup with $D(H) = 0$ has been identified, all the subgroups above H in the lattice must have $D(H) = 0$. Once again, there is no point in incorporating this into the computational algorithm.

A potentially useful extension of the calculation, beyond the scope of this paper, would be to construct the equations (1) that are consistent with the equivariance condition (2). This would permit automation of the explicit construction of solutions, which is of particular interest in cases where $D(H)$ is even, when solutions may or may not exist. This is a more difficult calculation, since it requires explicit construction of a matrix representation of G , which is non-unique, followed by the solution of a large number of equations, obtained from (2), relating the coefficients of nonlinear terms in (1).

4. Results

In this section, results are presented in tabular form, giving the isotropy subgroups for various finite groups of practical interest. Many of these results are new, although in some cases the results are available in the literature [11, 19, 12, 2]. In each case below, reference is made, where appropriate, to earlier work.

In some cases, general results can be obtained analytically. For the group D_n of symmetries of a regular n -gon, there are two-dimensional faithful irreps. If n is odd, there is one isotropy subgroup isomorphic to Z_2 , while if n is even there are two, corresponding to the one or two non-conjugate reflection symmetries [11].

Another example of an analytical result is available for the symmetric group S_n , where the natural n -dimensional representation is reducible and composed of the trivial irrep and an irrep of dimension $n - 1$. Here, the isotropy subgroups are of the form $S_p \times S_{n-p}$ for $p = 1 \dots [n/2]$, corresponding to a partition of n identical objects into two classes of size p and $n - p$ [10]. There is an analogous result for the alternating group A_n of even permutations. For $n \geq 4$, there is a natural irrep of dimension $n - 1$, with $[n/2]$ isotropy subgroups that are the stabilisers of the same partition. The subgroups are generated by even permutations of one subset and odd permutations of both subsets, so for $p = 1$ the axial isotropy subgroup is A_{n-1} , and for $p = 2$ it is isomorphic to S_{n-2} (since any permutation of $n - 2$ objects can be written as an even permutation of n objects, by interchanging the remaining two objects, or not, as required).

4.1. Symmetric, alternating and Platonic groups

In this section, results are presented in tabular form for some symmetric groups, alternating groups and the symmetry groups of the Platonic solids.

The results are given in the form of extended character tables for each group G . The first row of each table gives the number of elements in each conjugacy class, and the second row gives the orders of the elements in that class. Each subsequent row corresponds to a numbered irrep of G , marked with a symbol 'F' if the irrep is faithful, and gives the character for each class. This is followed by lists of isotropy subgroups with $D(H) = 1$ and $D(H) = 2$. The subgroups are labelled according to isotropy; in some cases, this does not completely define the subgroup, but to specify the group completely, by specifying its generators, would make the tables unwieldy. Larger groups that cannot easily be labelled are simply denoted H_n , where $n = |H|$. For isotropy subgroups with $N(H) \neq H$, the size of the quotient group $N(H)/H$ is given in brackets after the subgroup. Thus, if a group

with $D(H) = 1$ is followed by (2), the corresponding bifurcation must be a pitchfork. For subgroups with $D(H) = 2$, the subgroup is followed by the symbol ‘M’ if the subgroup is maximal (recall that all isotropy subgroups with $D(H) = 1$ are maximal).

For the alternating group A_4 (Table 2), which is also the symmetry group of rotations of a tetrahedron, there is one faithful 3D irrep, with axial subgroups Z_2 and Z_3 , corresponding to the two possible rotation axes of the tetrahedron.

The group A_5 is isomorphic to the rotation group of the icosahedron or dodecahedron (see Table 3). The problem of bifurcation with icosahedral symmetry has recently been investigated by Hoyle [12], who also computed the normal form equations for the bifurcation in each irrep, and derived stability conditions for each axial branch of solutions. For the 3D irreps, there are axial subgroups Z_2 , Z_3 and Z_5 , each corresponding to a rotation axis of the icosahedron, and in the natural 4D irrep there are axial subgroups $D_3 \equiv S_3$ and A_4 , in agreement with the general result for A_n given above. Note that in the 5D irrep, there is a maximal isotropy subgroup D_2 with $D(H) = 2$; it can be shown using the equations in [12] and the methods of [15, 16, 18] that solutions with isotropy D_2 exist, in addition to the D_3 and D_5 solutions guaranteed by the equivariant branching lemma.

For the group A_6 (Table 4), there are several cases where the equivariant branching lemma indicates branching to solutions with isomorphic but non-conjugate isotropy subgroups. In the natural 5D irrep, there are solutions with isotropy A_5 and S_4 , and an 18-element group generated by the permutations (1, 2, 3), (4, 5, 6) and (1, 2)(4, 5), as discussed above. In the 10D irrep, there are three maximal isotropy subgroups with $D(H) = 2$. Results have also been obtained for A_7 and A_8 (of order 20 160), where there are up to 21 axial subgroups.

Table 2: Results for the group A_4 , the group of rotations of a tetrahedron (see the beginning of Section 4.1 for notation).

	1	3	4	4		$D(H) = 1$	$D(H) = 2$
1	1	1	1	1	A_4		
2	1	1	$e^{2\pi i/3}$	$e^{-2\pi i/3}$	D_2		
3	1	1	$e^{-2\pi i/3}$	$e^{2\pi i/3}$	D_2		
4F	3	-1	0	0	Z_2 (2) Z_3		

Table 3: Results for the group A_5 , the group of rotations of an icosahedron; the symbols ϕ_{\pm} denote $(1 \pm \sqrt{5})/2$.

	1	15	20	12	12	$D(H) = 1$	$D(H) = 2$
1	1	1	1	1	1	A_5	
2F	3	-1	0	ϕ_+	ϕ_-	Z_2 (2) Z_3 (2) Z_5 (2)	
3F	3	-1	0	ϕ_-	ϕ_+	Z_2 (2) Z_3 (2) Z_5 (2)	
4F	4	0	1	-1	-1	D_3 A_4	Z_2 (2) Z_3 (2)
5F	5	1	-1	0	0	D_3 D_5	D_2 M (3)

Turning to the symmetric groups, Table 5 shows the case of S_4 , which has been discussed in earlier work [11, 19]. This group is isomorphic to the group of rotation symmetries of the cube (or octahedron), and also to the group of symmetries of the tetrahedron, including reflections. The ‘natural’ irrep for the rotations of the cube is that labelled 4F, and the axial subgroups correspond to rotations about the two-fold, three-fold and four-fold axes of the cube. Irrep 5F is the ‘natural’ one for the tetrahedron, or for the permutation of four objects. The isotropy subgroups D_2 and D_3 are analogous to the rotation axes found for A_4 , and the subgroup Z_2 with $D(H) = 2$ is generated by the reflection symmetry of the tetrahedron.

Table 4: Results for the group A_6 . The symbols ϕ_{\pm} denote $(1 \pm \sqrt{5})/2$.

	1	45	40	40	90	72	72
	1	2	3	3	4	5	5
1	1	1	1	1	1	1	1
2F	5	1	2	-1	-1	0	0
3F	5	1	-1	2	-1	0	0
4F	8	0	-1	-1	0	ϕ_+	ϕ_-
5F	8	0	-1	-1	0	ϕ_-	ϕ_+
6F	9	1	0	0	1	-1	-1
7F	10	-2	1	1	0	0	0

	$D(H) = 1$	$D(H) = 2$
1	A_6	
2F	$H_{18} (2) S_4 A_5$	$D_2 (6) D_3 A_4 (2)$
3F	$H_{18} (2) S_4 A_5$	$D_2 (6) D_3 A_4 (2)$
4F	$D_3 D_3 D_4 D_5$	$Z_3 (6) Z_3 (6) D_2 (6) D_2 (6) Z_4 (2) Z_5 (2)$
5F	$D_3 D_3 D_4 D_5$	$Z_3 (6) Z_3 (6) D_2 (6) D_2 (6) Z_4 (2) Z_5 (2)$
6F	$D_5 S_4 S_4 H_{36}$	$D_3 D_3 D_4$
7F	$D_3 D_3 A_4 (2) A_4 (2)$	$Z_4M (2) Z_5M (2) Z_3 \times Z_3M (4)$

Table 5: Results for the symmetric group S_4 .

	1	6	3	8	6		
	1	2	2	3	4	$D(H) = 1$	$D(H) = 2$
1	1	1	1	1	1	S_4	
2	1	-1	1	1	-1	$A_4 (2)$	
3	2	0	2	-1	0	D_4	$D_2 (6)$
4F	3	-1	-1	0	1	$Z_2 (2) Z_3 (2) Z_4 (2)$	
5F	3	1	-1	0	-1	$D_2 (2) D_3$	$Z_2 (2)$

For S_5 (see Table 6), the natural permutation irrep is 4F, giving axial subgroups S_4 and $S_3 \times S_2 \equiv D_6$. Results have also been found for S_6 , S_7 and S_8 .

The remaining symmetry groups for the Platonic solids are the groups $S_4 \times Z_2$ for the symmetry of the cube including reflections (Table 7), and $A_5 \times Z_2$ for the symmetry of the icosahedron with reflections (Table 8). The results are consistent with the earlier work [19, 12] in each case. For $S_4 \times Z_2$, the two faithful 3D irreps lead to axial subgroups D_2 , D_3 and D_4 . For $A_5 \times Z_2$, in the faithful 5D irrep there are axial subgroups D_2 , D_3 and D_5 ; as in the case of A_5 , there is a maximal two-dimensional isotropy subgroup D_2 . In this case, however, it can be shown (using the equations given in [12]) that there is generically no solution in this subspace.

Note that for a group G with a direct product structure $G = K \times Z_2$, the irreps can be obtained from those for the group K ; each irrep of K corresponds to two irreps ‘+’ and ‘-’ of G according to the action of the Z_2 subgroup. For the ‘+’ irreps, which are not faithful, the isotropy subgroups are simply the direct products of the isotropy subgroups for K with Z_2 , but for the faithful ‘-’ irreps there is no simple relationship between the isotropy subgroups of K and those of G .

Table 6: Results for the symmetric group S_5 .

	1	10	15	20	20	30	24
	1	2	2	3	6	4	5
1	1	1	1	1	1	1	1
2	1	-1	1	1	-1	-1	1
3F	4	-2	0	1	1	0	-1
4F	4	2	0	1	-1	0	-1
5F	5	-1	1	-1	-1	1	0
6F	5	1	1	-1	1	-1	0
7F	6	0	-2	0	0	0	1

	$D(H) = 1$	$D(H) = 2$
1	S_5	
2	$A_5 (2)$	
3F	$Z_4 (2) D_3 (2) Z_6 (2) A_4 (2)$	$Z_2 (4) Z_3 (4)$
4F	$D_6 S_4$	$D_2 (2) D_3 (2)$
5F	$D_3 (2) D_4 H_{20}$	$Z_2 (6) D_2 (6) Z_4 (2)$
6F	$D_4 D_5 (2) D_6$	$D_2 (6) D_2 (2)$
7F	$Z_4 (2) D_2 (2) D_3 (2) Z_6 (2)$	$Z_2 (4) Z_3 (4) Z_5M (4)$

Table 7: Results for the group $S_4 \times Z_2$ of symmetries of the cube including reflections.

	1	1	6	6	8	8	3	3	6	6
	1	2	2	2	3	6	2	2	4	4
1	1	1	1	1	1	1	1	1	1	1
2	1	1	-1	-1	1	1	1	1	-1	-1
3	1	-1	1	-1	1	-1	1	-1	1	-1
4	1	-1	-1	1	1	-1	1	-1	-1	1
5	2	2	0	0	-1	-1	2	2	0	0
6	2	-2	0	0	-1	1	2	-2	0	0
7F	3	-3	-1	1	0	0	-1	1	1	-1
8F	3	-3	1	-1	0	0	-1	1	-1	1
9	3	3	-1	-1	0	0	-1	-1	1	1
10	3	3	1	1	0	0	-1	-1	-1	-1

	$D(H) = 1$	$D(H) = 2$
1	H_{48}	
2	$H_{24} (2)$	
3	$S_4 (2)$	
4	$S_4 (2)$	
5	$D_4 \times Z_2$	$Z_2^3 (6)$
6	$D_4 (2) D_4 (2)$	$D_2 (12)$
7F	$D_2 (2) D_3 (2) D_4 (2)$	$Z_2 (4) Z_2 (8)$
8F	$D_2 (2) D_3 (2) D_4 (2)$	$Z_2 (4) Z_2 (8)$
9	$D_2 (2) Z_6 (2) Z_4 \times Z_2 (2)$	
10	$Z_2^3 (2) D_6$	$D_2 (2)$

Table 8: Results for the group $A_5 \times Z_2$ of symmetries of the icosahedron, including reflections.

	1	1	20	20	15	15	12	12	12	12
	1	2	3	6	2	2	5	10	5	10
1	1	1	1	1	1	1	1	1	1	1
2	1	-1	1	-1	1	-1	1	-1	1	-1
3	3	3	0	0	-1	-1	ϕ_+	ϕ_+	ϕ_-	ϕ_-
4F	3	-3	0	0	-1	1	ϕ_+	$-\phi_+$	ϕ_-	$-\phi_-$
5	3	3	0	0	-1	-1	ϕ_-	ϕ_-	ϕ_+	ϕ_+
6F	3	-3	0	0	-1	1	ϕ_-	$-\phi_-$	ϕ_+	$-\phi_+$
7	4	4	1	1	0	0	-1	-1	-1	-1
8F	4	-4	1	-1	0	0	-1	1	-1	1
9	5	5	-1	-1	1	1	0	0	0	0
10F	5	-5	-1	1	1	-1	0	0	0	0

	$D(H) = 1$	$D(H) = 2$
1	$A_5 \times Z_2$	
2	$A_5 (2)$	
3	$D_2 (2) Z_6 (2) Z_{10} (2)$	
4F	$D_2 (2) D_3 (2) D_5 (2)$	$Z_2 (4)$
5	$D_2 (2) Z_6 (2) Z_{10} (2)$	
6F	$D_2 (2) D_3 (2) D_5 (2)$	$Z_2 (4)$
7	$D_6 A_4 \times Z_2$	$D_2 (2) Z_6 (2)$
8F	$D_2 (2) D_3 (2) D_3 (2) A_4 (2)$	$Z_2 (4) Z_2 (4) Z_3 (4)$
9	$D_6 D_{10}$	$Z_2^3 M (3)$
10F	$D_2 (2) D_3 (2) D_5 (2)$	$Z_2 (4) D_2 M (6)$

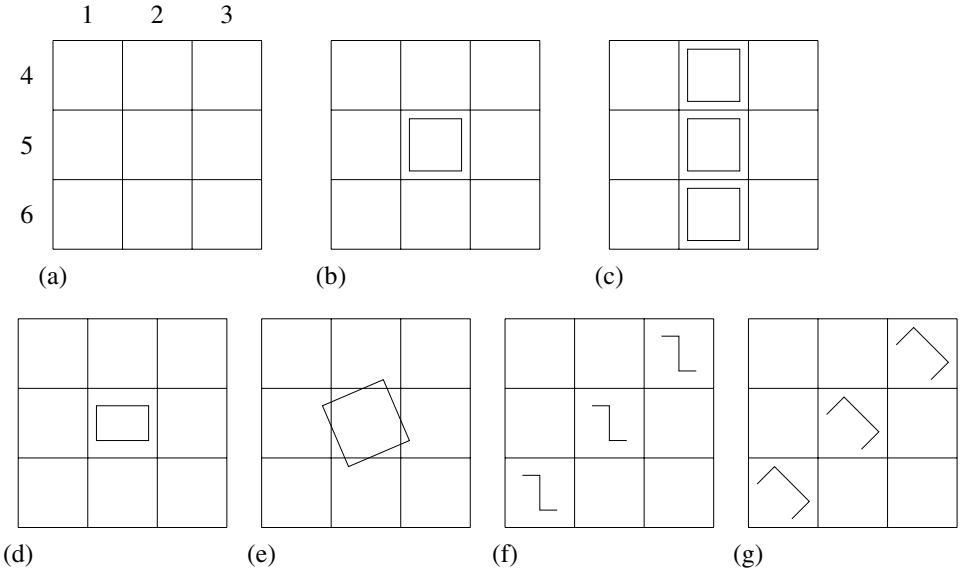


Figure 1: Symmetry-breaking from $D_4 \times Z_3 \times Z_3$.

- (a) The basic state and labelling used to write the group as a subgroup of S_6 .
- (b),(c) Schematic showing symmetry of axial solutions corresponding to irrep 6F with isotropy D_4 and D_6 .
- (d)–(g) Axial solutions corresponding to irrep 8F with isotropy D_2 , Z_4 , D_3 and Z_6 .

4.2. Instabilities of periodic patterns

An application of considerable current interest is the classification of instabilities of regular, spatially periodic patterns that break the periodicity of the pattern but maintain periodicity on some larger lengthscale. The resulting patterns are known as *superlattice* patterns. The examples discussed here are motivated by recent observations of transitions seen in the Faraday experiment [1, 14]. In this experiment, a layer of fluid or granular material is subjected to a vertical vibration. The uniform state can become unstable to a regular periodic pattern, which is often hexagonal. Also, by varying the parameters and introducing two-frequency forcing, a variety of secondary instabilities can be induced, leading to superlattice patterns. Since the underlying state in the Faraday experiment is periodic, the general theory of bifurcation from a stationary state outlined in Section 2 is not directly applicable. However, an analogous approach can be followed, in which the differential equations are replaced by maps describing the change in the state of the system after one cycle (see, for example, [24]).

Consider first a spatially periodic pattern of squares, with translation symmetries $x \rightarrow x + l$ and $y \rightarrow y + l$, and suppose that the symmetry of the pattern is broken in such a way that translation symmetries $x \rightarrow x + nl$ and $y \rightarrow y + nl$, $n \in \mathbb{N}$, are retained. With this constraint, the symmetry of the initial state is that of a periodic array of $n \times n$ squares, as shown in Figure 1(a) for $n = 3$. The symmetry group is generated by the symmetries D_4 of the square, together with the symmetries $Z_3 \times Z_3$ of the translations. Since the translations form a normal subgroup, the full symmetry group can be written as the semidirect product $G = D_4 \times Z_3 \times Z_3$. This group can conveniently be written

as a subgroup of the permutation group S_6 , by labelling the three rows and three columns as shown in Figure 1(a), and by writing down the permutation action of the group generators on the rows and columns. The reflection about a vertical axis interchanges columns 1 and 3, but fixes column 2 and the three rows, and so is represented by the permutation (1, 3). The reflection about the diagonal interchanges the rows and columns, corresponding to the permutation (1, 4)(2, 5)(3, 6). The translation $x \rightarrow x + l$ permutes the columns according to (1, 2, 3) (the translation $y \rightarrow y + l$ is not required to generate the group). The results of the symmetry-breaking computation are shown in Table 9. Since the subgroup $Z_3 \times Z_3$ is normal, the character table of the quotient group D_4 appears in the first five rows of the character table of G . These first five irreps correspond to instabilities of the square pattern that do not break any translation symmetries. There are four faithful four-dimensional irreps. In two of these, there are two axial subgroups, D_4 and D_6 . Further examination of the generators of these subgroups shows that in irrep 6F, the first pattern retains the square symmetry (1, 3), (1, 4)(2, 5)(3, 6) but no translation symmetry (shown schematically in Figure 1(b)), while the D_6 solution possesses a translation symmetry (4, 5, 6) and two orthogonal reflection symmetries (1, 3) and (4, 6) (see Figure 1(c)). In irrep 7F, the D_4 solution is the same, but the D_6 solution involves diagonal translations and reflections. Thus the D_6 subgroups in Table 9 are not conjugate, but the D_4 subgroups are conjugate. In these irreps there are no normaliser symmetries, so we expect these patterns to arise at a transcritical bifurcation. In the other four-dimensional irreps there are four competing patterns, with isotropies D_2 , Z_4 , D_3 and Z_6 . The solutions with D_2 and Z_4 symmetry have rectangular and $\pi/2$ rotation symmetry, respectively, and no translation symmetry (see Figure 1(d), (e)). The patterns with D_3 and Z_6 possess a translation symmetry combined with a π rotation symmetry and a reflection symmetry respectively (Figure 1(f), (g)). The Z_4 subgroups in irreps 8F and 9F are conjugate, but the D_2 , Z_6 and D_3 subgroups are not; they are related through a $\pi/4$ rotation. In these cases, the bifurcations are all pitchforks. In fact, the isotropy subgroups for $D_4 \times Z_n \times Z_n$ can be obtained analytically; the details of these calculations will be given elsewhere.

Table 9: Results for the group $D_4 \times Z_3 \times Z_3$ of symmetries of a 3×3 square lattice.

	1	2	3	2	6	3	2	4	6	$D(H) = 1$
1	1	1	1	1	1	1	1	1	1	$D_4 \times Z_3 \times Z_3$
2	1	-1	1	1	-1	1	1	-1	1	$D_2 \times Z_3 \times Z_3$ (2)
3	1	1	1	1	1	1	-1	-1	-1	$D_2 \times Z_3 \times Z_3$ (2)
4	1	-1	1	1	-1	1	-1	1	-1	$Z_4 \times Z_3 \times Z_3$ (2)
5	2	0	2	-2	0	2	0	0	0	$Z_2 \times Z_3 \times Z_3$ (2) $Z_2 \times Z_3 \times Z_3$ (2)
6F	4	2	1	0	-1	-2	0	0	0	$D_4 D_6$
7F	4	0	-2	0	0	1	2	0	-1	$D_4 D_6$
8F	4	-2	1	0	1	-2	0	0	0	D_2 (2) Z_4 (2) D_3 (2) Z_6 (2)
9F	4	0	-2	0	0	1	-2	0	1	D_2 (2) Z_4 (2) D_3 (2) Z_6 (2)

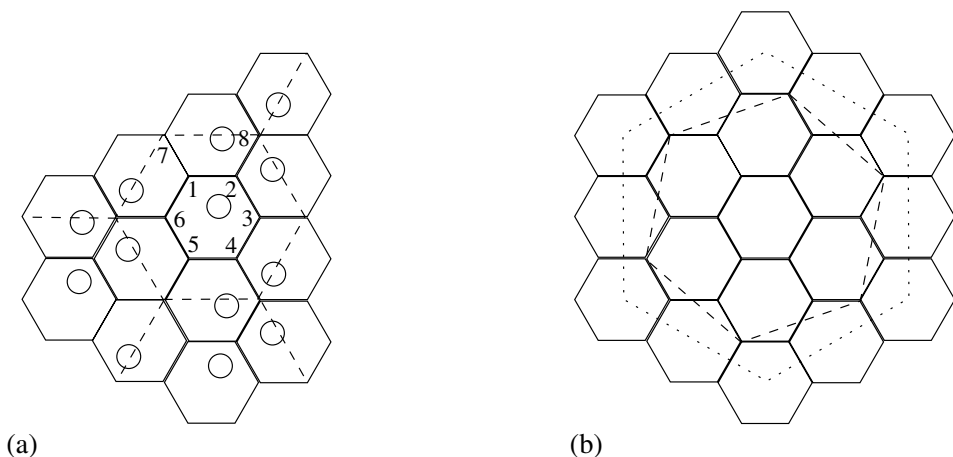


Figure 2: Symmetry-breaking on hexagonal lattices.

- (a) Circles indicate qualitatively the symmetry of the ‘SSS’ state of [1]. The pattern is periodic on a larger hexagonal lattice, indicated by dashed lines. The numbered vertices are used to embed the group in S_8 .
- (b) For the superlattice patterns observed in [14], the pattern is periodic on a larger hexagonal lattice containing seven hexagons (SL-I, dashed line) or twelve hexagons (SL-II, dotted line) [24].

In the experiments of Arbell and Fineberg [1], square patterns are found in some parameter regimes, and these can become unstable to a pattern that appears to have square symmetry with a threefold increase in the periodicity (states referred to as ‘2MS’, see [1, Figure 2(b)]). The symmetry of these states is approximately as shown in Figure 1(b), so that the instability corresponds to the 4D irrep 6F or 7F. There are axial solutions with D_4 and D_6 symmetry, but in the experiment it appears that only the D_4 solution is stable.

Arbell and Fineberg also observed hexagonal states, and these can become unstable to a state in which the hexagons cluster in groups of four (‘SSS’ states, see [1, Figure 2(a)]). This transition can be analysed by considering a hexagonal array set within a larger hexagonal lattice, so that each large hexagon contains four small hexagons (Figure 2(a)). In this case, the relevant symmetry group is $D_6 \times Z_2 \times Z_2$, of order 48, generated by the hexagonal symmetry and the two translations of the lattice. To write the group as a permutation group, the eight distinct vertices of the hexagons are labelled as shown in Figure 2. The generators are then the reflections $(1, 2)(3, 6)(4, 5)(7, 8)$ and $(1, 3)(4, 6)$ and the translation $(1, 8)(2, 6)(3, 5)(4, 7)$. Investigating the structure of the group reveals that it is isomorphic to the group of symmetries of the cube, $S_4 \times Z_2$, the Z_2 subgroup being generated by a rotation through π about the centre of a hexagon. Hence the possible symmetry-breaking instabilities are exactly analogous to those of the cube, and the axial subgroups can be read off from Table 7. The pattern observed in [1] and sketched in Figure 2(a) has D_2 symmetry generated by the reflections $(1, 3)(4, 6)$ and $(1, 4)(2, 8)(3, 6)(5, 7)$. The product of these reflections is a π rotation about the edge of a hexagon. This subgroup is axial in the three-dimensional faithful irreps 7F and 8F in Table 7. (Note that the pattern does not have the D_2 symmetry of irreps 9 and 10, since in these irreps the π rotation about the centre of a hexagon acts as the identity, so all the patterns have that symmetry.)

In the Faraday experiments of Kudrolli, Pier and Gollub [14], two types of superlattice patterns were observed. The first type (SL-I) is periodic on a skewed hexagonal lattice containing seven hexagons, and the second (SL-II) is periodic on a lattice containing twelve hexagons (Figure 2(b)). The SL-II patterns were studied using the group-theoretic approach by Tse *et al.* [24]. In this case, the symmetry group can be embedded in S_{12} simply by labelling the twelve hexagons and writing down the action of the two reflection symmetries and a translation symmetry. The translations generate the group $Z_6 \times Z_2$, so the full spatial symmetry group of the system is $D_6 \times Z_6 \times Z_2$, of order 144 (see [24]). The character table is shown in Table 10 (note that the character table given in [24] contains minor typographical

Table 10: Character table for the group $G = D_6 \times Z_6 \times Z_2$ of symmetries of a periodic lattice of 12 hexagons, with axial subgroups for G in irreps 14F and 15F and axial subgroups for $G \times Z_2$ in irrep 15F–.

	1	6	18	24	8	3	9	16	12	12	6	18	2	6	3
	1	2	2	6	3	2	2	3	12	6	6	4	3	4	2
1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
2	1	1	-1	-1	1	-1	-1	1	1	1	1	-1	1	1	1
3	1	-1	-1	1	1	1	1	1	-1	-1	1	-1	1	-1	1
4	1	-1	1	-1	1	-1	-1	1	-1	-1	1	1	1	-1	1
5	2	0	0	-1	-1	2	2	-1	0	0	2	0	2	0	2
6	2	0	0	1	-1	-2	-2	-1	0	0	2	0	2	0	2
7	2	2	0	0	2	0	0	-1	-1	-1	-1	0	-1	2	2
8	2	-2	0	0	2	0	0	-1	1	1	-1	0	-1	-2	2
9	3	-1	-1	0	0	3	-1	0	1	-1	-1	1	3	1	-1
10	3	-1	1	0	0	-3	1	0	1	-1	-1	-1	3	1	-1
11	3	1	-1	0	0	-3	1	0	-1	1	-1	1	3	-1	-1
12	3	1	1	0	0	3	-1	0	-1	1	-1	-1	3	-1	-1
13	4	0	0	0	-2	0	0	1	0	0	-2	0	-2	0	4
14F	6	-2	0	0	0	0	0	0	-1	1	1	0	-3	2	-2
15F	6	2	0	0	0	0	0	0	1	-1	1	0	-3	-2	-2

14F	$D_2 (2) D_2 (2) Z_6 (2) D_3 (2) D_3 D_4 (2) Z_4 \times Z_2 (2)$
15F	$D_3 D_4 (2) Z_2^3 (2) D_6$
15F–	$D_3 (2) D_3 (2) Z_2^3 (2) Z_2^3 (2) D_6 (2) D_6 (2) D_4 \times Z_2 (2) D_4 \times Z_2 (2)$

errors that do not affect the results). In fact, the group is isomorphic to $S_4 \times S_3$, although this is by no means obvious and does not generalise to other superlattices. Note that there are two faithful six-dimensional irreps. It is shown in [24] that the relevant irrep for the SL-II pattern of [14] is 15F. The axial subgroups for the faithful irreps 14F and 15F are listed in Table 10; there are four axial subgroups for irrep 15F.

An additional complication in this example is that the observed instability is subharmonic; that is, it arises at a period-doubling bifurcation, with Floquet multiplier -1 . It is explained in [24] that this can be incorporated by including an additional Z_2 symmetry to represent the time translation, so that the group is now isomorphic to $(D_6 \times Z_6 \times Z_2) \times Z_2$, of order 288. As discussed at the end of Section 4.1, each irrep of the original group then splits into two irreps, ‘+’ and ‘-’, according to the action of the additional Z_2 symmetry. The irrep of interest is then the 15F- irrep, and the axial subgroups for this irrep are listed in the last row of Table 10. Tse *et al.* [24] state that there are six axial isotropy subgroups for this irrep, but do not explain how they did the calculation. Since there are 272 subgroups, it is likely that the axial subgroups were found by inspection, rather than by following the systematic algorithm outlined in Table 1. In fact, the computational algorithm reveals that there are eight axial subgroups. The pair of subgroups isomorphic to D_3 are omitted in [24]; the solution corresponding to the SL-II pattern of [14] is identified in [24] as one of the solutions with isotropy $D_4 \times Z_2$. Most of the axial subgroups are ‘twisted’; that is, they include spatiotemporal symmetries involving a spatial symmetry and a time shift. However, one of the D_3 solutions and one of the D_6 solutions correspond to states with purely spatial symmetries. These states have triangular symmetry about a vertex of one of the original hexagons, and hexagonal symmetry about the centre of one of the original hexagons, respectively.

5. Oscillatory bifurcations

The preceding sections have been concerned entirely with the case of a stationary bifurcation, where a real eigenvalue passes through zero, leading to branches of stationary solutions. We now consider the case of an oscillatory, or Hopf, bifurcation, where a complex conjugate pair of eigenvalues crosses into the right half-plane and the bifurcating branches are periodic orbits.

For the case of a Hopf bifurcation, it is necessary to consider symmetry-breaking from the group $G \times S^1$, where S^1 denotes the continuous time symmetry that is broken in the Hopf bifurcation (see [9, 10, 11] for details of the general theory). The resulting solution branches have spatiotemporal symmetries, involving combinations of elements of G with time translations. More precisely, the spatiotemporal symmetry H contains a normal subgroup K that is an isotropy subgroup of G , and H/K is a cyclic group representing the time translations. This cyclic group H/K must be maximally abelian in $N(K)/K$. These conditions are necessary for the Hopf bifurcation form of the equivariant branching lemma to be valid [9]. For the dihedral group D_n , in the two-dimensional irreps, it is well known that there are three solution branches [11]. The subgroups H (with spatial subgroup K in brackets) are $Z_n(1)$, $Z_2(1)$, $Z_2(Z_2)$ for n odd, and $Z_n(1)$, $D_2(Z_2)$, $D_2(Z_2)$ for n even.

Since the time translation symmetry is a cyclic group, a convenient algorithm for reducing the case of the Hopf bifurcation to the case of the stationary bifurcation is to consider the finite group $G \times Z_m$, for some choice of m . It is straightforward to see that the time shift symmetry H/K must have an order that is equal to the order of one of the elements of G , so m can be chosen as the least common multiple of the orders of the elements of G .

For example, $m = 4$ for D_4 , $m = 12$ for S_4 , and $m = 30$ for A_5 . In the cases when m is large, an alternative approach is to consider the different possible values of m separately; for example, for A_5 , the calculation is repeated for $A_5 \times Z_2$, $A_5 \times Z_3$ and $A_5 \times Z_5$.

The results obtained using this algorithm are listed in Table 11. For each group, only the faithful irreps are considered, and these are listed using the same labelling as in the tables in Section 4 for the stationary bifurcation. For each solution, the full spatiotemporal symmetry H is given, as well as the normal subgroup K representing the spatial part. In some cases $H = K$, indicating that the symmetries of the periodic orbit are purely spatial.

In a few cases, the results in Table 11 had been obtained previously. Swift and Barany [23] considered Hopf bifurcation with the symmetry of rotation of a tetrahedron (A_4). They found the four branches of periodic solutions shown in Table 11, and showed that the normal form equations can exhibit chaotic behaviour. For the group S_4 , which is the group of rotations of a cube, Ashwin and Podvigina [2] studied the Hopf bifurcation in irrep 4F; their results are in agreement with Table 11. In irrep 5F, the axial solutions have isotropy that is isomorphic to those for irrep 4F. This is because the characters for these two irreps differ only in signs (Table 5), and these signs can be changed by a half shift in time. Thus, although the spatiotemporal symmetry groups are isomorphic, the spatial subgroups are not.

Table 11: Axial solution branches for Hopf bifurcation. Labels for the irreps are as used on earlier tables for the stationary bifurcation. The spatiotemporal symmetry for each axial solution is given, with the spatial subgroup K in brackets.

Group	Irrep	Axial branches
A_4	4F	$Z_3 (1) Z_3 (1) Z_3 (Z_3) D_2 (Z_2)$
A_5	2F	$Z_3 (1) Z_5 (1) D_2 (Z_2) D_3 (Z_3) D_5 (Z_5)$
A_5	3F	$Z_3 (1) Z_5 (1) D_2 (Z_2) D_3 (Z_3) D_5 (Z_5)$
A_5	4F	$Z_3 (1) Z_5 (1) Z_5 (1) D_2 (Z_2) D_3 (Z_3) D_3 (D_3) A_4 (A_4)$
A_5	5F	$Z_5 (1) Z_5 (1) D_2 (Z_2) D_3 (D_3) D_5 (D_5) A_4 (D_2) A_4 (D_2)$
S_4	4F	$Z_3 (1) Z_4 (1) D_2 (Z_2) D_3 (Z_3) D_4 (Z_4)$
S_4	5F	$Z_3 (1) Z_4 (1) D_2 (Z_2) D_3 (D_3) D_4 (D_2)$
$S_4 \times Z_2$	7F	$Z_2^3 (D_2) Z_6 (1) Z_4 \times Z_2 (Z_2) D_6 (D_3) D_4 \times Z_2 (D_4)$
$S_4 \times Z_2$	8F	$Z_2^3 (D_2) Z_6 (1) Z_4 \times Z_2 (Z_2) D_6 (D_3) D_4 \times Z_2 (D_4)$
$A_5 \times Z_2$	4F	$Z_2^3 (D_2) D_6 (D_3) D_{10} (D_5) Z_6 (1) Z_{10} (1)$
$A_5 \times Z_2$	6F	$Z_2^3 (D_2) D_6 (D_3) D_{10} (D_5) Z_6 (1) Z_{10} (1)$
$A_5 \times Z_2$	8F	$Z_2^3 (D_2) D_6 (D_3) D_6 (D_3) A_4 \times Z_2 (A_4) Z_6 (1) Z_{10} (1) Z_{10} (1)$
$A_5 \times Z_2$	10F	$Z_2^3 (D_2) D_6 (D_3) D_{10} (D_5) A_4 \times Z_2 (D_2) A_4 \times Z_2 (D_2) Z_{10} (1) Z_{10} (1)$

For the Hopf bifurcation with the symmetry of the cube including reflections, there are again five branches, and the two faithful irreps 7F and 8F of Table 7 give the same results. For Hopf bifurcation with icosahedral symmetry, there are seven distinct branches of periodic solutions in the four- and five-dimensional irreps. The calculation was also carried out for the group $D_4 \times Z_3 \times Z_3$ (Figure 1), corresponding to an oscillatory instability of a regular square pattern. In this case, there are nine solution branches for each of the faithful four-dimensional irreps of Table 9.

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References

1. H. ARBELL and J. FINEBERG, ‘Spatial and temporal dynamics of two interacting modes in parametrically driven surface waves’, *Phys. Rev. Lett.* 81 (1998) 4384–4387. 102, 112, 114
2. P. ASHWIN and O. PODVIGINA, ‘Hopf bifurcation with cubic symmetry and instability of ABC flow’, *Proc. Roy. Soc. A* 459 (2003) 1801–1827. 106, 117
3. E. BODENSCHATZ, W. PESCH and G. AHLERS, ‘Recent developments in Rayleigh–Bénard convection’, *Ann. Rev. Fluid. Mech.* 32 (2000) 709–778. 102
4. P. CHOSSAT and R. LAUTERBACH, *Methods in equivariant bifurcations and dynamical systems*, Adv. Ser. Nonlin. Dynam. 15 (World Scientific, River Edge, NJ, 2000). 101, 102, 103
5. G. CICOGNA, ‘Symmetry breakdown from bifurcations’, *Lettere al Nuovo Cimento* 31 (1981) 600–602. 103
6. S. M. COX and P. C. MATTHEWS, ‘Instability and localisation of patterns due to a conserved quantity’, *Physica D* 175 (2003) 196–219. 102
7. J. H. P. DAWES, P. C. MATTHEWS and A. M. RUCKLIDGE, ‘Reducible actions of $D_4 \times T^2$: superlattice patterns and hidden symmetries’, *Nonlinearity* 16 (2003) 615–645. 102, 103
8. The GAP Group, GAP – *Groups, algorithms, and programming*, Version 4.3 (2002); <http://www.gap-system.org>. 102, 104
9. M. GOLUBITSKY and I. STEWART, ‘An algebraic criterion for symmetric Hopf bifurcation’, *Proc. Roy. Soc. London Ser. A* 440 (1993) 727–732. 116
10. M. GOLUBITSKY and I. STEWART, *The symmetry perspective* (Birkhäuser, Basel, 2002). 101, 103, 106, 116
11. M. GOLUBITSKY, I. STEWART and D. G. SCHAEFFER, *Singularities and groups in bifurcation theory*, vol. II, Appl. Math. Ser. 69 (Springer, New York, 1988). 101, 102, 103, 105, 106, 108, 116
12. R. B. HOYLE, ‘Shapes and cycles arising at the steady bifurcation with icosahedral symmetry’, *Physica D* 191 (2004) 261–281. 106, 107, 109
13. E. IHRIG and M. GOLUBITSKY, ‘Pattern selection with $O(3)$ symmetry’, *Physica D* 13 (1984) 1–33. 103, 105
14. A. KUDROLLI, B. PIER and J. P. GOLLUB, ‘Superlattice patterns in surface waves’, *Physica D* 123 (1998) 99–111. 102, 112, 114, 115, 116

15. A. LARI-LAVASSANI, W. F. LANGFORD and K. HUSEYIN, ‘Symmetry-breaking bifurcations on multidimensional fixed point subspaces’, *Dynam. Stability Systems* 9 (1994) 345–373. [103](#), [107](#)
16. R. LAUTERBACH, ‘Spontaneous symmetry breaking in higher-dimensional fixed point spaces’, *Z. Angew. Math. Phys.* 43 (1992) 430–448. [103](#), [107](#)
17. M. J. LINEHAN and G. E. STEDMAN, ‘Little groups of irreps of $O(3)$, $SO(3)$, and the infinite axial subgroups’, *J. Phys. A* 34 (2001) 6663–6688. [105](#)
18. P. C. MATTHEWS, ‘Bifurcation in two-dimensional fixed point subspaces’, preprint (2003); arXiv:math.DS/0305392. [103](#), [107](#)
19. I. MELBOURNE, ‘A singularity theory analysis of bifurcation problems with octahedral symmetry’, *Dynamics and Stability of Systems* 1 (1986) 293–321. [106](#), [108](#), [109](#)
20. M. MICHEL, ‘Symmetry defects and broken symmetry’, *Rev. Modern Phys.* 52 (1980) 617–651. [105](#)
21. LORD RAYLEIGH, ‘On convection currents in a horizontal layer of fluid, when the higher temperature is on the underside’, *Philos. Mag.* 32 (1916) 529–546. [102](#)
22. D. H. SATTINGER, *Group theoretic methods in bifurcation theory* (Springer, New York, 1979). [101](#)
23. J. W. SWIFT and E. BARANY, ‘Chaos in the Hopf bifurcation with tetrahedral symmetry: convection in a rotating fluid with low Prandtl number’, *Proc., Nonlinear Hydrodynamic Stability and Transition, Sophia-Antipolis, 1990, European J. Mech. B Fluids* 10 (1991) 99–104. [117](#)
24. D. P. TSE, A. M. RUCKLIDGE, R. B. HOYLE and M. SILBER, ‘Spatial period-multiplying instabilities of hexagonal Faraday waves’, *Physica D* 146 (2000) 367–387. [102](#), [112](#), [114](#), [115](#), [116](#)
25. A. VANDERBAUWHEDE, ‘Symmetry and bifurcation near families of solution’, *J. Differential Equations* 36 (1980) 173–178. [103](#)

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